

Combined constraints on dark photons from high-energy collisions, cosmology, and astrophysics

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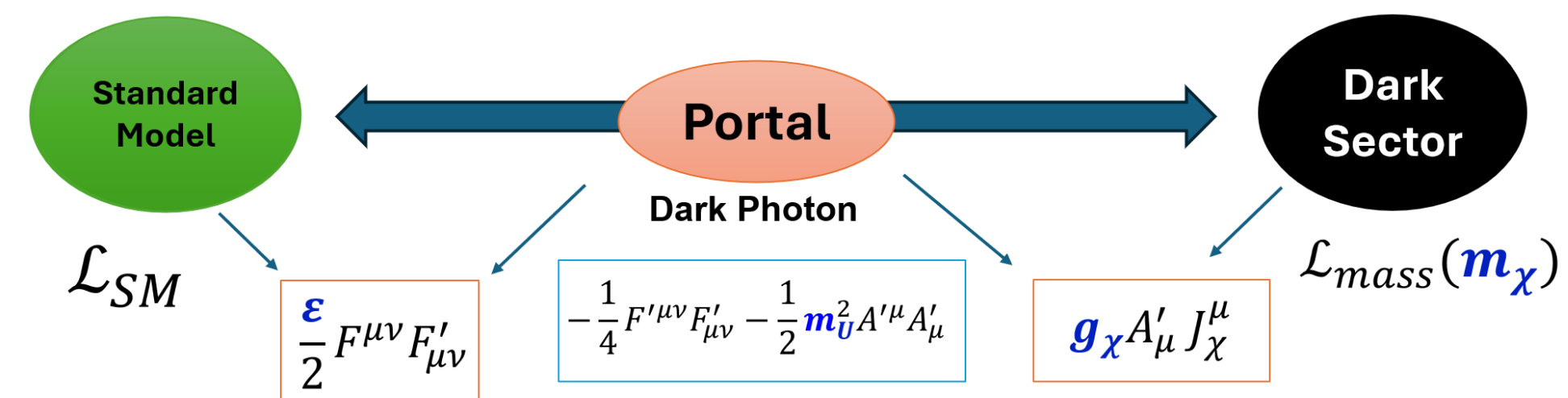
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1. Motivation

- Goal: Constrain vector portal dark photons by combining dilepton spectra with dark matter (DM) relic density and self-interacting DM (SIDM) constraints.

2. Vector portal with Dark Matter

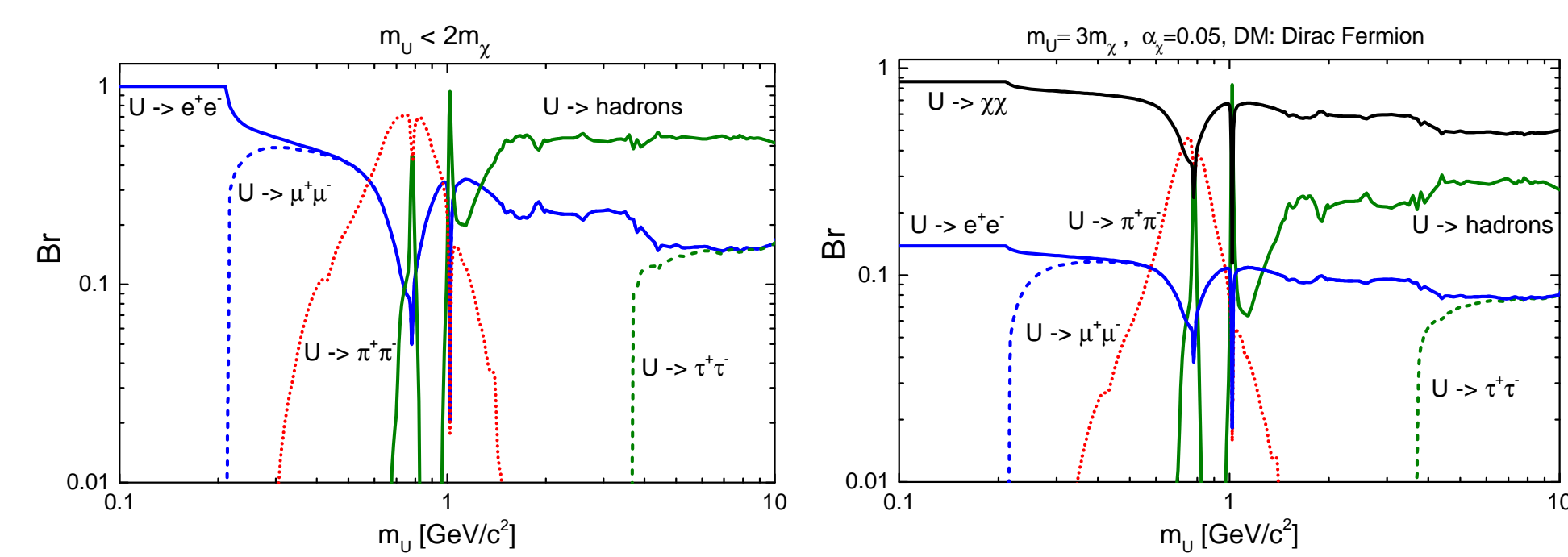
The dark photon lagrangian with dark matter is [1]



At low energies, the mediator couples as $\epsilon e A'_\mu J_{EM}^\mu$. The dark current J_χ^μ reads as:

$$J_\chi^\mu = \begin{cases} \bar{\chi} \gamma^\mu \chi, & \text{Dirac fermion,} \\ \frac{1}{2} \bar{\chi} \gamma^\mu \gamma^5 \chi, & \text{Majorana fermion,} \\ i(\varphi^\dagger \partial^\mu \varphi - (\partial^\mu \varphi^\dagger) \varphi). & \text{complex scalar,} \end{cases}$$

Visible ($m_U < 2m_\chi$) and Invisible ($m_U > 2m_\chi$) width



- For $m_U < 2m_\chi$, the dilepton branching ratio ($U \rightarrow e^+e^-$) remains unsuppressed. For $m_U > 2m_\chi$, invisible decays reduce the observable $U \rightarrow e^+e^-$ signal.

3. Dark photon production in PHSD

The Parton-Hadron-String Dynamics (PHSD) is a non-equilibrium off-shell transport approach for relativistic hadronic and heavy-ion collisions. It describes the full space-time evolution of the collision, from the initial nucleon-nucleon interactions and string excitation, through the QGP phase, to hadronization and final-state hadronic rescattering and decays. [3]



Dark-photon production channels in PHSD [2]

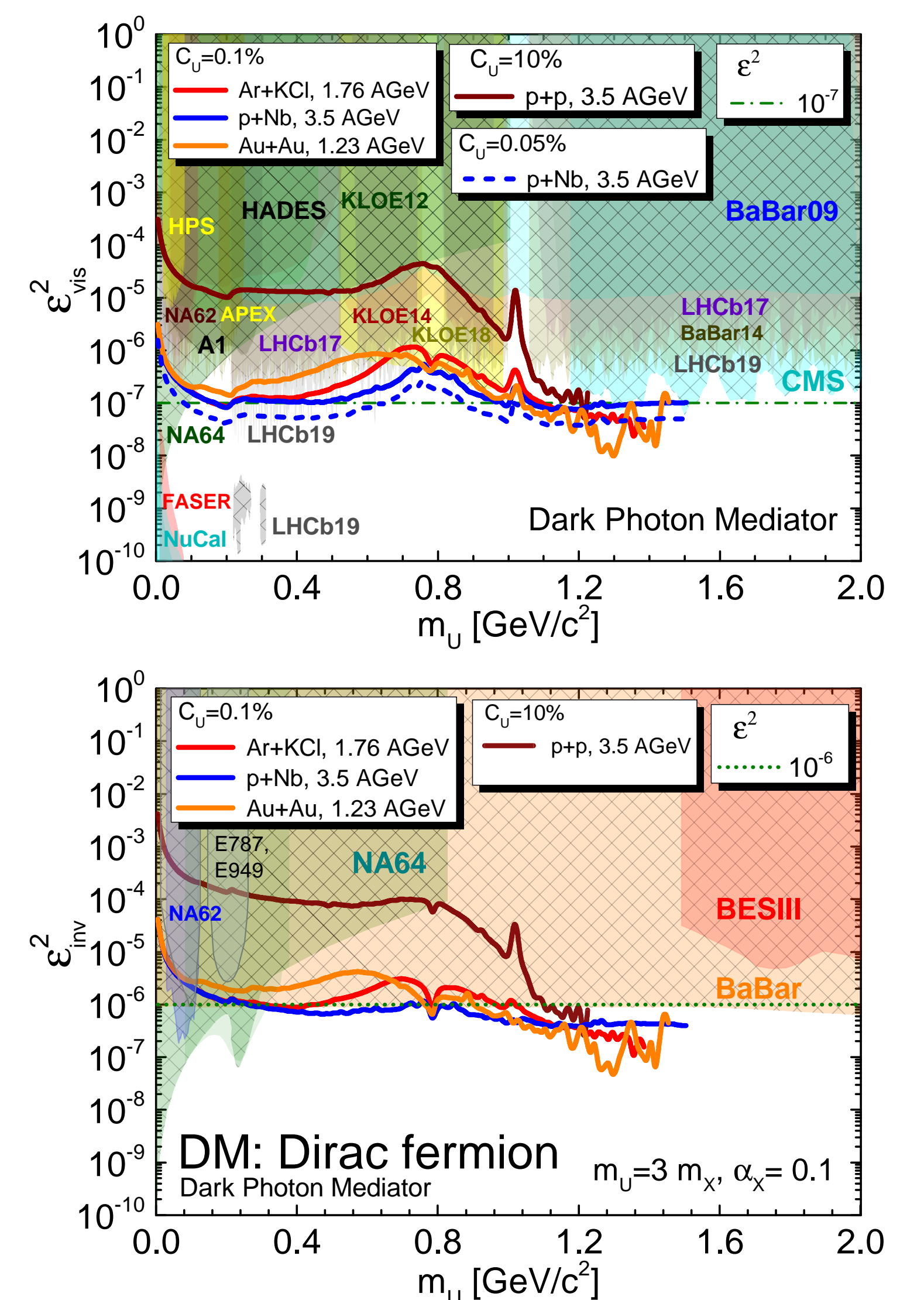
$$\begin{aligned} \pi^0, \eta, \eta' &\rightarrow \gamma U, & \rho, \omega, \phi &\rightarrow U, \\ \omega &\rightarrow \pi^0 U, & K^+ &\rightarrow \pi^+ U, \\ \Delta &\rightarrow N U, & q\bar{q} &\rightarrow U. \end{aligned}$$

The total dilepton yield from dark photons is

$$N^{U \rightarrow e^+e^-} = \text{Br}(U \rightarrow e^+e^-) \sum_h N_{h \rightarrow UX}$$

- Hadronic sources are implemented explicitly in PHSD [2].
- The partonic contribution is obtained by rescaling the PHSD quark-gluon plasma dilepton yield.

4. From dilepton excess to a limit on ϵ^2



For each invariant-mass bin dM , the dark-photon yield scales linearly with kinetic mixing square,

$$\frac{dN^{\text{sum}U}}{dM} = \epsilon^2 \frac{dN^{\text{sum}U}}{dM}.$$

We require the dark photon contribution not to exceed a small surplus over the SM yield,

$$\frac{dN^{\text{sum}U}}{dM} \leq C_U \frac{dN^{\text{sum}SM}}{dM}.$$

This yields the upper bound for the kinetic mixing [2]

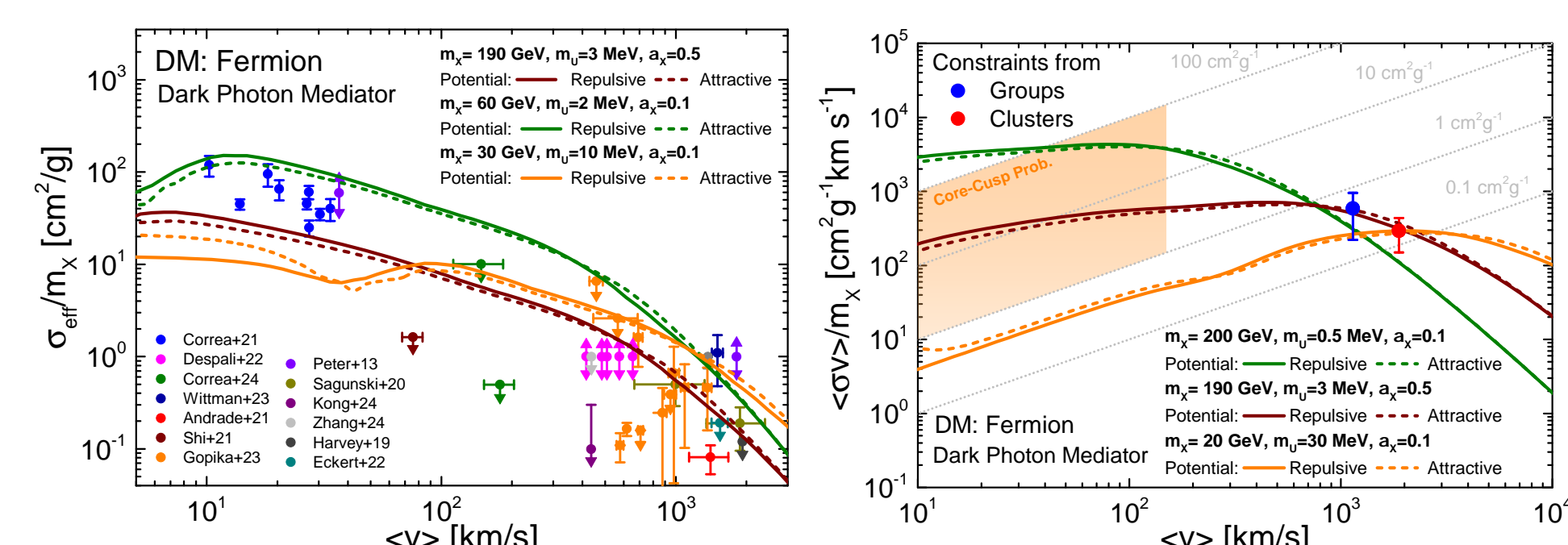
$$\epsilon^2(m_U) = C_U \left(\frac{dN^{\text{sum}SM}/dM}{dN^{\text{sum}U}/dM} \right)$$

- Heavy-ion dilepton data probe both visible and invisible dark-photon regimes.

6. Self-interacting dark matter constraints

We compute Yukawa-mediated DM self-scattering with CLASSICS, and use transport cross-sections [4, 6]

$$\sigma_T(v) = \int d\Omega (1 - \cos\theta) \frac{d\sigma}{d\Omega}, \quad \sigma_{\text{eff}} = \frac{3 \langle \sigma(v) v^5 \rangle}{2 \langle v^5 \rangle}.$$

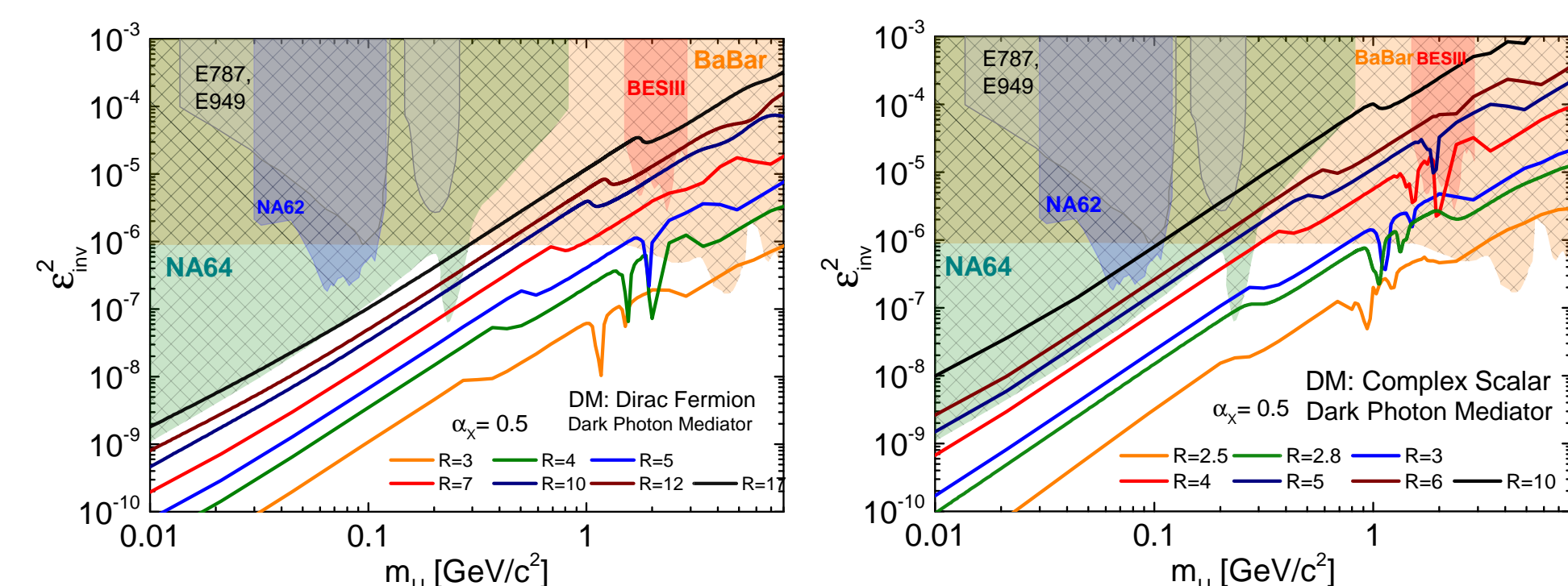


- Viable SIDM scenarios require large self-interactions at dwarf velocities and suppressed scattering at group and cluster scales.
- This typically favors light mediators and comparatively heavy dark matter.

7. Thermal relic target

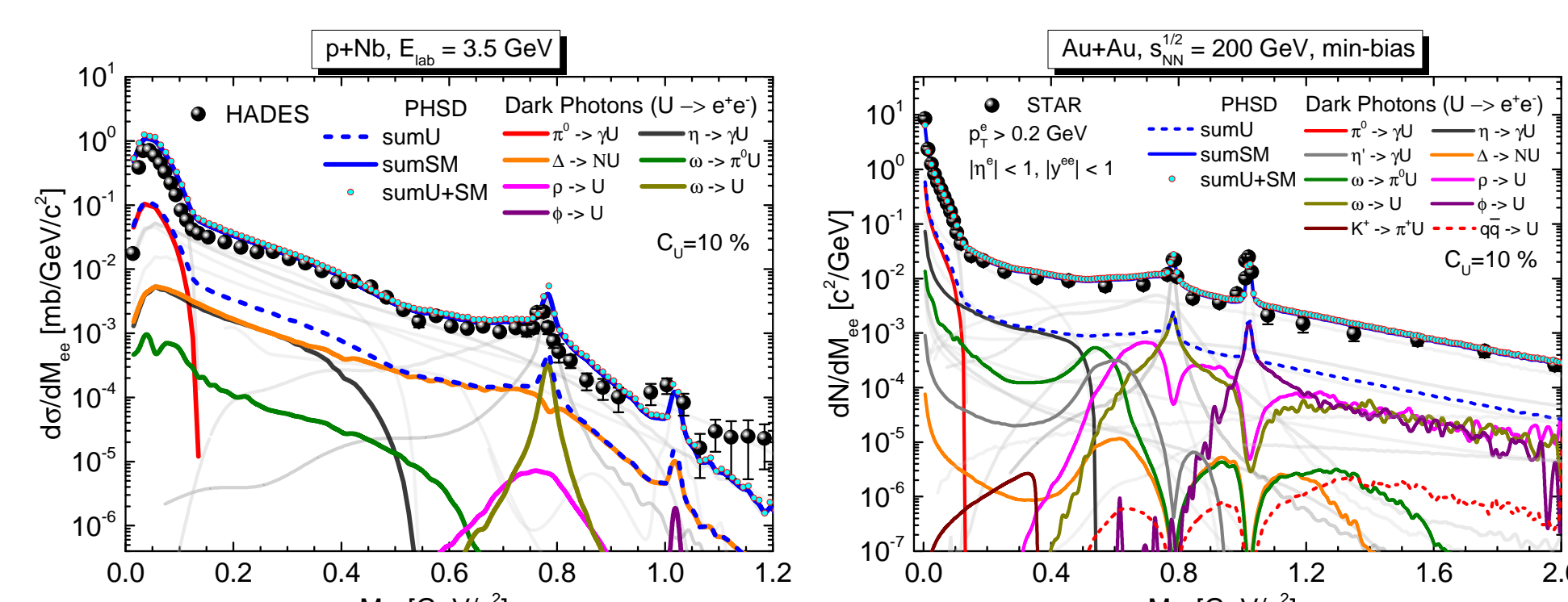
Using RED-DELIVER, we solve the freeze-out dynamics and require the observed DM relic abundance [5],

$$\Omega_{\text{DM}} h^2 \simeq 0.120 \pm 0.001.$$



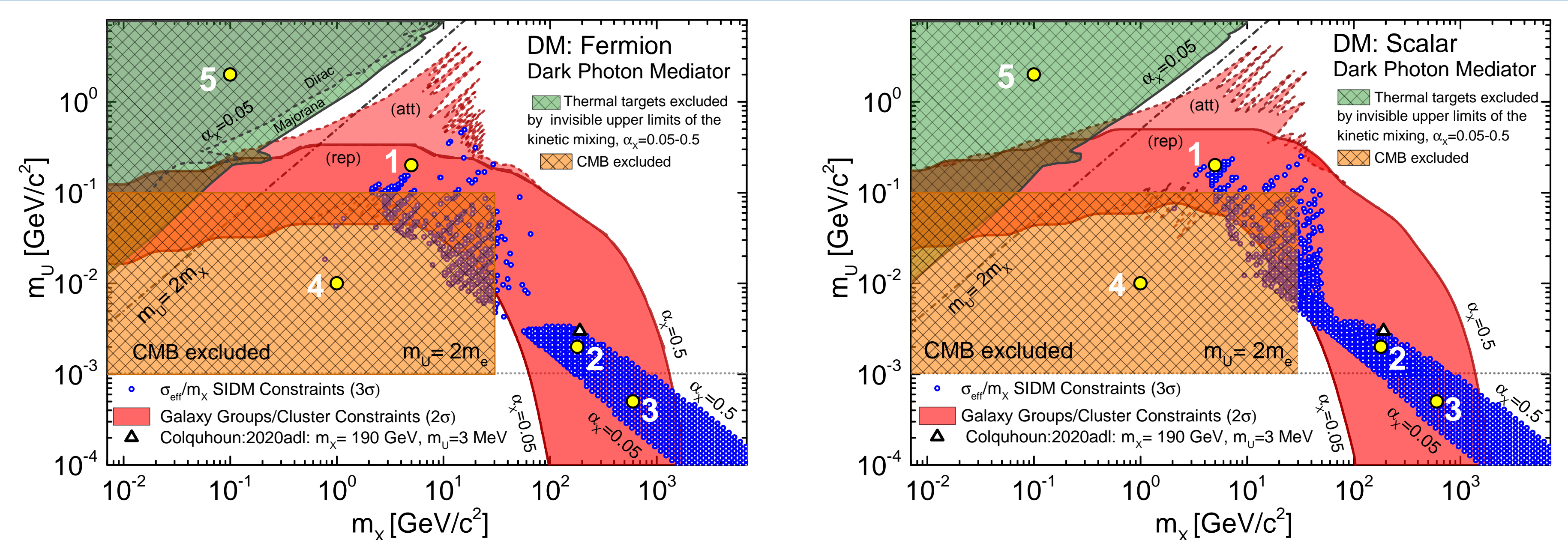
- Dirac, Majorana, and complex scalar DM lead to distinct thermal target curves.
- Local dips originate from hadronic thresholds and resonance effects in the mediator width.

5. PHSD dilepton spectra with dark photons



- PHSD reproduces the Standard Model dilepton background from SIS to LHC energies. [1–3]
- Dark photons appear as narrow structures in the invariant-mass spectra.
- Improved experimental precision translates directly into stronger limits on ϵ^2 through smaller C_U .

8. Combined parameter-space constraints



- Blue/red:** Regions compatible with SIDM (effective cross-section and Groups/Cluster constraints) from dwarfs, galaxies, groups, and clusters.
- Orange:** CMB-excluded DM low-mass region.
- Green:** Thermal target points excluded by invisible kinetic mixing limits on ϵ^2 .
- $m_U = 2m_\chi$ Separates the visible and invisible regimes.

- BP1:** visible sub-GeV mediator and intermediate-mass DM, it satisfies SIDM constraints.
- BP2:** MeV mediator with heavy DM, typical of SIDM-favored scenarios.
- BP3:** Ultra-light long-lived mediator with $m_U < 2m_e$: stable dark photon.
- BP4/BP5:** Representative CMB and PHSD excluded benchmarks (see orange and green regions).

9. Summary

- We have studied a dark sector in which a kinetically mixed dark photon U is constrained through dilepton production within the PHSD transport approach.
- Combining the kinetic mixing limits with thermal relic targets and self-interaction constraints, we found that large regions of the parameter space are excluded.
- The combined analysis favors light mediators in the MeV-sub-GeV range and comparatively heavy dark matter (DM),

KEY REFERENCES

- [1] A. W. Romero Jorge, L. Sagunski, et al., Phys. Rev. D 113 (2026) 055052. [2] A. W. Romero Jorge, E. Bratkovskaya, et al., Phys. Rev. C 112 (2025) 054905. [3] P. Moreau, O. Solovleva et al., Phys. Rev. C 100 (2019) 014911. [4] B. Colquhoun, S. Heeba et al., Phys. Rev. D 103 (2021) 035006. [5] N. Aghanim et al. (Planck), Astron. Astrophys. 641 (2020) A6. [6] M. S. Fischer, L. Kasselmann, et al., Mon. Not. Roy. Astron. Soc. 529 (2024) 2327.

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